Precise Measurement of the CP Violation Parameter \( \sin 2\phi_1 \) in \( B^0 \to (c\bar{c})K^0 \) Decays


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In the standard model (SM), CP violation in the quark sector is described by the Kobayashi-Maskawa (KM) theory [1] in which the quark-mixing matrix has a single irreducible complex phase that gives rise to all CP-violating asymmetries. In the decay chain $Y(4S) \rightarrow B^0 \bar{B}^0 \rightarrow f_C P f_{\text{tag}}$, where one of the $B$ mesons decays at time $t_{\text{CP}}$ to a $CP$ eigenstate $f_{CP}$, and the other $B$ meson decays at time $t_{\text{tag}}$ to a final state $f_{\text{tag}}$ that distinguishes between $B^0$ and $\bar{B}^0$, the decay rate has a time dependence in the $Y(4S)$ rest frame [2] given by

$$P(\Delta t) = \frac{e^{-|\Delta t|/\tau_{B^0}}}{4\tau_{B^0}} \left[ 1 + q[S_f \sin(\Delta m_d \Delta t) + A_f \cos(\Delta m_d \Delta t)] \right].$$

(1)

Here $S_f$ and $A_f$ are $CP$ violation parameters, $\tau_{B^0}$ is the $B^0$ lifetime, $\Delta m_d$ is the mass difference between the two neutral $B$ mass eigenstates, $\Delta t \equiv t_{\text{CP}} - t_{\text{tag}}$, and the $b$-flavor charge $q = +1(-1)$ when the tagging $B$ meson is a $B^0 (\bar{B}^0)$. With very small theoretical uncertainty [2], the SM predicts $S_f = -\xi_f \sin2\phi_1$ and $A_f = 0$ for the $b \rightarrow c\bar{c}s$ transition, where $\xi_f = +1(-1)$ corresponds to $CP$-even (-odd) final states and $\phi_1$ is an interior angle of the KM unitarity triangle, defined as $\phi_1 \equiv \text{arg}[-V_{cd}V_{cb}^*/V_{td}V_{tb}^*]$ [3]. The BABAR and Belle Collaborations have published several determinations of $\sin2\phi_1$ since the first observation [4,5]; previous results used $465 \times 10^6$ [6] and $535 \times 10^6$ [7] $BB$ pairs, respectively.

With recently available experimental results, not only $\sin2\phi_1$ but also other measurements of the sides of the unitarity triangle and other $CP$ violation measurements make it possible to test the consistency of the KM scheme. The indirect determination of the angle $\phi_1$ deviates by $2.7\sigma$ from the current world average for the direct determination of $\sin2\phi_1$ [8]. Equivalently, the $B^\pm \rightarrow \tau^\pm \nu_\tau$ branching fraction and the resulting value of $|V_{ub}|$ differ by $2.8\sigma$ from the prediction of the global fit [8], where the
sin2ϕ1 value gives the most stringent constraint on the indirect measurement. Furthermore, time-dependent CP violation in the neutral B meson decays mediated by flavor-changing b → s transitions may deviate from CP violation in the b → c̄es case because of possible additional quantum loops [9]. To clarify whether new physics contributes to CP-violating phenomena or B^± → τ^±ντ decays, it is very important to determine sin2ϕ1, the SM reference, as precisely as possible.

In this Letter, we describe the final Belle measurement of sin2ϕ1 and δϕ1 in b → c̄es induced B decays to fCP. The B decays to the CP-odd eigenstates, fCP = J/ψK^0_S, ψ(2S)K^0_S, and χc1K^0_S, and the CP-even eigenstate, fCP = J/ψK^0_L, are reconstructed using 772 × 10^6 BB pairs, the entire data sample accumulated on the Y(4S) resonance with the Belle detector [10] at the KEKB asymmetric-energy e^+e^− collider [11]. Two inner detector configurations were used. A 2.0 cm radius beam pipe and a 3-layer silicon vertex detector (SVD) were used for the first data sample that contains 152 × 10^6 BB pairs. The remaining 620 × 10^6 BB pairs were accumulated with a 1.5 cm radius beam pipe, a 4-layer silicon vertex detector, and a small-cell inner drift chamber. The latter data sample has been recently reprocessed using a new charged track reconstruction algorithm, which significantly increased the reconstruction efficiency for the B^0 → (c̄e)K^0_S decay modes. In particular, the gain for the B^0 → J/ψK^0_S decay mode is 18%.

The Y(4S) is produced with a Lorentz boost of βγ = 0.425 nearly along the z axis, which is antiparallel to the positron beam direction. Since the B^0 and B^± mesons are approximately at rest in the Y(4S) center-of-mass (c.m.) system, Δt can be determined from the displacement in z between the fCP and f_μg decay vertices: Δt ≈ (zCP − z_μg)/(βγc) ≡ Δz/(βγc).

Charged tracks reconstructed in the central drift chamber (CDC), except for tracks from K^0_S → π^+π^− decays, are required to originate from the interaction point (IP). We distinguish charged kaons from pions based on a kaon (pion) likelihood L_K(π) derived from the time-of-flight scintillation counters, aerogel threshold Cherenkov counters, and dE/dx measurements in the CDC. Electron identification is based on the ratio of the electromagnetic calorimeter (ECL) cluster energy to the particle momentum as well as a combination of dE/dx measurements in the CDC, the aerogel threshold Cherenkov response, and the position and shape of the electromagnetic shower. Muons are identified by track penetration depth and hit scatter in the muon detector (KLM). Photons are identified as isolated ECL clusters that are not matched to any charged track.

For the J/ψK^0_S, J/ψK^0_L, and ψ(2S)K^0_S modes, event selection is the same as in our previous analyses [7,12], where J/ψ mesons are reconstructed via their decays to ℓ^+ℓ^− (ℓ = e, μ) and the ψ(2S) mesons to ℓ^+ℓ^− or J/ψπ^+π^−. For the modes J/ψK^0_L and χc1K^0_S, in which the χc1 is reconstructed in the J/ψγ final state, both J/ψ daughter tracks must be positively identified as leptons, whereas for the J/ψK^0_S and ψ(2S)K^0_S modes, at least one daughter must satisfy this requirement. Any other track having an ECL energy deposit consistent with a minimum ionization particle is accepted as a muon candidate and any track satisfying either the dE/dx or the ECL shower energy requirements is retained as an electron candidate. For J/ψ → e^+e^− decays, the e^± e^± decay daughters are combined with photons found within 50 mrad of the e^+ or e^− direction in order to account partially for final-state radiation and bremsstrahlung. In order to accommodate the remaining radiative tails, an asymmetric invariant mass requirement is used to select J/ψ and ψ(2S) decays in dilepton modes, −150 MeV/c^2 < M_{e^+e^-} < M_{μ^+μ^-} < 36 MeV/c^2 and −60 MeV/c^2 < M_{e^+e^-} < 36 MeV/c^2, where M_{μ^+μ^-} denotes either the nominal J/ψ or ψ(2S) mass. For ψ(2S) → J/ψπ^+π^− candidates, we require a mass difference of 580 MeV/c^2 < M_{e^+e^-} < 600 MeV/c^2, and χc1 → J/ψγ candidates are required to have a mass difference of 385.0 MeV/c^2 < M_{e^+e^-} < 430.5 MeV/c^2. For each charmonium candidate, vertex-constrained and mass-constrained fits are applied to improve its momentum resolution.

Candidate K^0_S → π^+π^− decays are selected by requirements on their invariant mass, flight length, and consistency between the K^0_S momentum direction and vertex position. Candidate K^0_S mesons are selected from ECL and/or KLM hit patterns that are consistent with the presence of a shower induced by a K^0_L meson. The centroid of the K^0_S candidate shower is required to be within a 45° cone centered on the K^0_S direction calculated from the two-body B decay kinematics and the momentum of the reconstructed J/ψ meson.

For B → fCP candidate reconstruction in modes other than J/ψK^0_L, B candidates are identified by two kinematic variables: the energy difference ΔE = E^*_B − E^*_beam and the beam-energy constrained mass M_{bc} ≡ √(E^*_beam)^2 − (p^*_B)^2, where E^*_beam is the CM beam energy, and E^*_beam(p^*_B) is the CM energy (momentum) of the reconstructed B candidate. The B^0 → J/ψK^0_L candidates are identified by the value of p^*_B calculated using a two-body decay kinematic assumption.

The b flavor of the accompanying B meson is identified from inclusive properties of particles that are not associated with the reconstructed B^0 → f_CP decay [13]. The tagging information is represented by two parameters, the b-flavor charge q and purity r. The parameter r is an event-by-event, MC-determined flavor-tagging dilution factor that ranges from r = 0 for no flavor discrimination to r = 1 for unambiguous flavor assignment. The data are sorted into seven intervals of r. For events with r > 0.1, the wrong-tag fractions for six r intervals, w_j (j = 1, 6), and their differences between B^0 and B^0 decays, Δw_j, are...
determined from semileptonic and hadronic $b \rightarrow c$ decays [12,14]. If \( r \approx 0.1 \), the wrong-tag fraction is set to 0.5, and therefore the tagging information is not used. The total effective tagging efficiency, \( \Sigma(f_i \times (1 - 2w_i))^2 \), is determined to be \( 0.298 \pm 0.004 \), where \( f_i \) is the fraction of events in the category \( i \).

The vertex position for the \( f_{CP} \) decay is reconstructed using \( J/\psi \) or \( \psi(2S) \) daughter tracks that have a minimum number of SVD hits, while the \( f_{tag} \) vertex is determined from well-reconstructed tracks that are not assigned to \( f_{CP} \) [14]. A constraint on the IP profile in the plane perpendicular to the \( z \) axis is used with the selected tracks. With this procedure, we are able to determine a vertex even in the case where only one track has sufficient associated SVD hits. The fractions of the single track vertices for \( f_{CP} \) and \( f_{tag} \) are about 12% and 23%, respectively.

For a single track vertex, the estimated error of the \( z \) coordinate, \( \sigma_z \), is the indicator of the vertex fit quality and is required to be less than 500 \( \mu \)m. On the other hand, a vertex reconstructed using two or more tracks is characterized by a more robust goodness-of-fit indicator. In the previous analysis [7], the value of \( \chi^2 \) of the vertex calculated solely along the \( z \) direction was used. This is now replaced by \( h \), the value of \( \chi^2 \) in three-dimensional space calculated using the charged tracks \textit{without} using the interaction-region profile’s constraint [15]. A detailed MC study indicates that \( h \) is a superior indicator of the vertex goodness-of-fit because it is less sensitive to the specific \( B \) decay mode; in particular, \( h \) shows a smaller mode dependence for the vertices reconstructed from \( B \rightarrow J/\psi X \) and \( B \rightarrow D^{(*)}X \) decays, which are used as control samples to determine the vertex resolution parameters. In the multiple-track vertex case, \( h < 50 \) and \( \sigma_z < 200 \mu \)m are required. For candidate events in which both \( B \) vertices are reconstructed, we retain only those events where the \( B \) vertices satisfy \( |\Delta t| < 70 \) ps for further analysis.

For the candidate events in which both flavor tagging and vertex reconstruction succeed, the signal yield and purity for each mode are obtained from an unbinned maximum-likelihood fit to the two-dimensional \( \Delta E - M_{bc} \) distribution for \( f_{CP} \) modes with a \( K^0_L \) meson, and to the \( p_B^* \) distribution for \( J/\psi K^0_L \). The background mainly comes from \( B\bar{B} \) events in which one of the \( B \) mesons decays into a final state containing a correctly reconstructed \( J/\psi \), i.e., the \( B \rightarrow J/\psi X \) process. In order to determine this background distribution, a \( B \rightarrow J/\psi X \) MC sample corresponding to 100 times the integrated luminosity of data is used. An estimate of other combinatorial background is obtained from the \( M_{F->} \) sideband. For \( CP \)-odd modes, the signal distribution is modeled with a Gaussian function in \( M_{bc} \) and a double Gaussian function in \( \Delta E \). The fits to determine signal yields for these modes are performed in the region \( 5.2 \) GeV/\( c^2 < M_{bc} < 5.3 \) GeV/\( c^2 \) and \( -0.1 \) GeV < \( \Delta E < 0.2 \) GeV. The \( p_B^* \) signal shape for \( J/\psi K^0_L \) is determined from MC events. The requirement \( p_B^* < 2.0 \) GeV/c is used in the fit to estimate the signal yield as well as the contribution of three categories of background: those with a real (those that are correctly reconstructed) \( J/\psi \) and a real \( K^0_L \), those with a real \( J/\psi \) and a fake \( K^0_L \) (those that are incorrectly reconstructed from electronic noise or electromagnetic showers), and events with a fake \( J/\psi \) (those that are background combinations). The \( M_{bc} \) distribution for a stringent \( \Delta E \) requirement (\( |\Delta E| < 40 \) MeV for \( J/\psi K^0_L \), \( |\Delta E| < 30 \) MeV for \( \psi(2S)K^0_L \), and \( |\Delta E| < 25 \) MeV for \( \chi_{c1}K^0_L \)) as well as the \( p_B^* \) distribution for \( J/\psi K^0_L \) candidates is shown in Fig. 1. We require \( 5.27 \) GeV/c/\( c^2 < M_{bc} < 5.29 \) GeV/c/\( c^2 \) for \( f_{CP} \) modes with a \( K^0_L \) and 0.20 GeV/c < \( p_B^* < 0.45 \) GeV/c for \( J/\psi K^0_L \) for the fit to the \( CP \) violation parameters. For the candidates passing all the criteria mentioned above, the signal yield and purity are estimated for each \( CP \) eigenstate and are listed in Table 1.

We determine \( S_f \) and \( \mathcal{A}_f \) for each mode by performing an unbinned maximum-likelihood fit to the observed \( \Delta t \) distribution. The probability density function (PDF) for the signal distribution, \( \mathcal{P}_{signal}(\Delta t; S_f, \mathcal{A}_f, q, w_f, \Delta \omega_f) \), is given by Eq. (1), fixing \( \tau_{B^0} \) and \( \Delta m_{\psi} \) at their world average values [16] and including modifications to take the effect of an incorrect flavor assignment (parametrized by \( w_f \) and \( \Delta \omega_f \)) into account. The distribution is convolved with the proper-time interval resolution function, \( R_{sig}(\Delta t) \), formed by convolving four components: the detector resolutions for \( z_{CP} \) and \( z_{tag} \), the shift of the \( z_{tag} \) vertex position due to secondary tracks from charmed particle decays, and the kinematic approximation that the \( B \) mesons are at rest in the CM frame [17]. Because we now use \( h \) to characterize the vertex goodness of fit, each of these resolution function components in Ref. [17] is reformulated as a function of \( h \) and \( \sigma_z \).

Using the \( M_{bc} \) sideband events, the background PDF, \( \mathcal{P}_{bkg}(\Delta t) \), for each of the \( CP \)-odd modes is modeled as a

\[ \mathcal{P}_{signal}(\Delta t; S_f, \mathcal{A}_f, q, w_f, \Delta \omega_f) \]

\[ \mathcal{P}_{bkg}(\Delta t) \]
sum of exponential and prompt components, and is convolved with $R_{bg}(\Delta t)$ expressed as a double Gaussian function. In the $J/\psi K_S^0$ mode, there are CP violating modes among the $B \to J/\psi X$ backgrounds, which are included in the background PDF. The $\Delta t$ PDFs for the remaining $B \to J/\psi X$ and other combinatorial backgrounds are estimated from the corresponding large MC sample and $M_{c\bar{c}}^{-}$ - sideband events, respectively. The construction of these PDFs follows the same procedure as in our previous analyses [7,12].

We determine the following likelihood for the $i$th event:

$$P_i = (1 - f_{ol}) \sum_k f_k \int \left[ P_k(\Delta t') R_k(\Delta t_i - \Delta t') \right] d(\Delta t')$$

$$+ f_{ol} P_{ol}(\Delta t_i),$$

(2)

where the index $k$ labels each signal or background component. The fraction $f_k$ depends on the $r$ region and is calculated on an event-by-event basis as a function of $\Delta E$ and $M_{bc}$ for the CP-odd modes and $P^F_r$ for the CP-even mode. The term $P_{ol}(\Delta t)$ is a broad Gaussian function that represents an outlier component $f_{ol}$, which has a fractional normalization of order 0.5% [17]. The only free parameters in the fits are $S_j$ and $A_j$, which are determined by maximizing the likelihood function $L = \prod_i P_i(\Delta t_i; S_j, A_j)$. This likelihood is maximized for each $f_{CP}$ mode individually, as well as for all modes combined taking into account their CP-eigenstate values; the results are shown in Table II. Figure 2 shows the $\Delta t$ distributions and asymmetries for good tag quality ($r > 0.5$) events. We define the background-subtracted asymmetry in each $\Delta t$ bin by

$$\left( N_++N_- \right) / \left( N_+ + N_- \right),$$

where $N_+$ ($N_-$) is the signal yield with $q = +1 (-1)$.

Uncertainties originating from the vertex reconstruction algorithm are a significant part of the systematic error for both $\sin 2 \phi_1$ and $A_{CP}$. These uncertainties are reduced by almost a factor of 2 compared to the previous analysis [7] by using $h$ for the vertex-reconstruction goodness-of-fit parameter, as described above. In particular, the effect of the vertex quality cut is estimated by changing the requirement on either $h < 25$ or $h < 100$; the systematic error due to the IP constraint in the vertex reconstruction is estimated by varying the IP profile size in the plane perpendicular to the $z$ axis; the effect of the criterion for the selection of tracks used in the $f_{tag}$ vertex is estimated by changing the requirement on the distance of the closest approach with respect to the reconstructed vertex by $\pm 50 \mu m$ from the nominal maximum value of 500 $\mu m$. Systematic errors due to imperfect SVD alignment are estimated from MC samples that have artificial misalignment effects. Small biases in the $\Delta z$ measurement are observed in $e^+ e^- \to \mu^+ \mu^-$ and other control samples. To account for these, a special correction function is applied and the variation with respect to the nominal results is included as a systematic error. We also vary the $|\Delta t|$ range by $\pm 30$ ps to estimate the systematic uncertainty due to the $|\Delta t|$ fit range. The vertex resolution function is another major source of $\sin 2 \phi_1$ and $A_{CP}$ uncertainty. This effect is estimated by varying each resolution function parameter obtained from data (MC) by $\pm 1\sigma$ ($\pm 2\sigma$) and repeating the fit to add each variation in quadrature. The uncertainty in the estimated errors of the parameters of reconstructed charged tracks is also taken into account. The largest contribution to the

<table>
<thead>
<tr>
<th>Decay mode</th>
<th>$\xi_f$</th>
<th>$N_{sig}$</th>
<th>Purity (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>$J/\psi K_S^0$</td>
<td>$-1$</td>
<td>$12649 \pm 114$</td>
<td>97</td>
</tr>
<tr>
<td>$\psi(2S)(\ell^+\ell^-)K_S^0$</td>
<td>$-1$</td>
<td>$904 \pm 31$</td>
<td>92</td>
</tr>
<tr>
<td>$\psi(2S)(J/\psi \pi^+\pi^-)K_S^0$</td>
<td>$-1$</td>
<td>$1067 \pm 33$</td>
<td>90</td>
</tr>
<tr>
<td>$\chi_{c1}K_S^0$</td>
<td>$-1$</td>
<td>$940 \pm 33$</td>
<td>86</td>
</tr>
<tr>
<td>$J/\psi K_L^0$</td>
<td>$+1$</td>
<td>$10040 \pm 154$</td>
<td>63</td>
</tr>
</tbody>
</table>

TABLE I. $CP$ eigenvalue ($\xi_f$), signal yield ($N_{sig}$), and purity for each $B^0 \to f_{CP}$ mode.
systematic uncertainty in $A_f$ is the effect of the tag-side interference (TSI), which is described in detail in [18]. Since the effect of TSI has an opposite sign for different CP eigenstates, there is a partial cancellation in the combined result. Hence, the combined TSI systematic is smaller than the systematic in each individual mode. Systematic errors due to uncertainties in the wrong-tag fractions are studied by varying the wrong-tag fraction individually in each $r$ region. A possible fit bias is examined by fitting a large number of MC events. Other contributions come from uncertainties in the signal fractions, the background $\Delta r$ distribution, $\tau_{pi}$, and $\Delta m_d$. Each contribution is summarized in Table III. We add them in quadrature to obtain the total systematic uncertainty.

In summary, we present the final $\sin 2 \phi_1$ measurement using the entire Belle Y(4S) data sample containing 772 $\times$ 10$^6$ $B\bar{B}$ pairs. We have reconstructed $b \rightarrow c \bar{s} s$ induced $B$ meson decays in three CP-odd modes [$J/\psi K^0_S$, $\psi(2S)K^0_S$, and $\chi_{c1} K^0_S$] and one CP-even mode ($J/\psi K^0_L$). The fit, using common CP-sensitive parameters for all four modes, yields the values $\sin 2 \phi_1 = 0.667 \pm 0.023$ (stat) $\pm 0.012$ (syst) and $A_f = 0.006 \pm 0.016$ (stat) $\pm 0.012$ (syst). The results are consistent with previous measurements [6,7]. These are the most precise determinations of these parameters and solidify the SM reference value used to test for evidence of new physics beyond the SM.

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[3] Another naming convention $\beta (= \phi_1)$ is also used in the literature.